



Research Article

Structural, electronic, and optical properties of N/Y-codoped anatase TiO₂: A first-principles investigation

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Abstract

The influence of N/Y codoping on the structural, electronic, and optical characteristics of anatase TiO₂ is investigated using density function theory. The calculated band gaps are 2.149 eV for pure TiO₂, 1.534 eV for N-doped TiO₂, 2.142 eV for Y-doped TiO₂, and 1.880 eV for N/Y-codoped TiO₂. The slightly lower value of pure TiO₂ compared to the experimental band gap (~3.2 eV) is attributed to the well-known underestimation effect of the GGA-PBESol functional in DFT calculations. The valence band of pure TiO₂ is mainly composed of O 2*p* states, and the conduction band is mainly composed of Ti 3*d* states. For N-doped TiO₂, the valence band is mainly occupied by N 2*p* states, which are higher than that of the O 2*p* states, which can narrow the band gap. Y doping will introduce 4*d* states, which are involved in the conduction band. As for N/Y codoped, the valence band is mainly occupied by N 2*p* and O 2*p* states, while the conduction band is predominantly occupied by Y 4*d* and Ti 3*d* states. Furthermore, N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ will result in a red-shift of the absorption edge compared with the pure TiO₂ and significantly enhanced in the visible light region.

Keywords: First-principles calculations, density functional theory, TiO₂, N/Y-codoped TiO₂, electronic properties, optical properties

Introduction

Many research articles say that the Honda-Fujishima effect is the first observation of the photocatalytic effect published in Nature in 1972 [1]. A previous work on photocatalytic water splitting using TiO₂ photoelectrodes by Fujishima and Honda marked the dawn of a new era in solar hydrogen harvesting using artificial photosynthesis. In earlier days, TiO₂ was used as a white pigment because it is a low-cost, light-stable, and environmentally friendly material. TiO₂ was mixed with clothing in the 20th century, and it was used for degradation. The photobleaching of dyes by TiO₂ was first reported by Doodeve and Kitchener in 1938 [2]. In addition, the degradation mechanism also explains the UV-absorbed TiO₂, creating the surface-reactive oxygen species known as photosensitizers. Mashio et al. subsequently published on "TiO₂ photocatalyst auto-oxidation" in 1956 [3]. It has been found that H₂O₂ is produced

during the photocatalytic reaction. The photocatalytic work was exponentially improved after the discovery of these remarkable results, which were evident from recent scientific publications. Photovoltaic water electrolysis may become more competitive as the cost continues to decrease with technological advancement. However, the considerable use of small band gap semiconducting materials may cause serious life cycle environmental impacts.

Most of the semiconductor photocatalytic materials, such as gallium arsenide (GaAs), lead sulfide (PbS), cadmium sulfide (CdS), zinc oxide (ZnO), iron oxide (Fe₂O₃), tin oxide (SnO₂), and tungsten trioxide (WO₃), have been examined as semiconductor-based photocatalysts [4–7]. Among them, TiO₂ is much more promising as it is stable, non-corrosive, environmentally friendly, abundant, and cost-effective. Rutile, anatase, and brookite are the three

naturally occurring forms of TiO₂. However, because of the complex structure of brookite and difficulties synthesising it in its purest form, more studies are focused on rutile and anatase TiO₂. Anatase TiO₂ is more favourable than rutile and brookite as it has better performance for photocatalytic hydrogen production due to faster charge carrier separation and higher kinetic stability [8, 9]. Unfortunately, the wide band gap of anatase TiO₂ (~3.2 eV) gives a limitation for its light usage in the ultraviolet (UV) region. Therefore, one of the successful methods to reduce the band gap of TiO₂ and thus increase the light harvesting in photocatalytic applications is by doping TiO₂ with non-metal and metal elements.

Doping TiO₂ is an important strategy to modify the band gap and optical response of semiconductor photocatalysts. The main objective of doping is to induce a bathochromic shift as a result decrease in the band gap or introduction of intra-band gap states, which leads to the absorption of more visible light. Doping may lead to photocatalytic systems that exhibit enhanced efficiency [10]. Many studies have been conducted to investigate the doping effect of N-doped TiO₂ theoretically and experimentally [11–16]. Doping by non-metal ions could change the electronic structure and improve the visible light absorption of the photocatalyst by means of narrowing its band gap [17–19]; however, the narrower band gap results in an increased recombination rate of photo-excited carriers, reducing the photocatalytic efficiency. Thus, transition metal element doping can effectively reduce the electron-hole recombination rate. Various research groups have studied TiO₂ doped with many different transition metals [20–24]. Recent studies suggested that transition metal and non-metal codoped in TiO₂ can effectively extend their optical absorption edges from the ultraviolet region to the visible light region and prevent the recombination of electron-hole pairs [25–29].

In this work, the structural, electronic, and optical properties of N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ were studied using the exchange–correlation functional from the GGA method based on the density functional theory (DFT). The doping of N into TiO₂ can reduce its band gap and shift the optical absorption edge toward the visible light region, while the doping of Y into TiO₂ can effectively reduce the electron-hole recombination rate. Both effects are beneficial for improving the performance of the N/Y-codoped TiO₂ photocatalyst. The various manufacturing methods and experimental settings have complex effects on the photocatalytic capabilities of TiO₂; hence, the mechanism of doping

is still unclear. Therefore, N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ using DFT need to be studied to clarify the effects on host structural, electronic, and optical properties.

Computational details

Anatase is a metastable mineral form of titanium dioxide (TiO₂). The tetragonal structure of anatase TiO₂ belongs to the spatial group I41/*amd* with lattice parameters $a = b = 3.720 \text{ \AA}$ and $c = 9.600 \text{ \AA}$. The non-metal and metal-doped anatase TiO₂ systems using N and Y as dopants were constructed by replacing one Ti atom with a Y atom and one O atom with an N atom in the anatase supercell. In the N/Y-codoped configuration, the N atom substituted an O site, and its adjacent Ti atom was replaced by a Y atom, as shown in Fig. 1. This nearest-neighbor substitution model was selected because it has been reported to be one of the most energetically favorable and structurally stable configurations for N-metal codoped TiO₂ systems. Such close dopant proximity facilitates effective charge transfer and orbital hybridization between N 2*p* and Y 4*d* states, which enhances visible-light absorption. Additionally, the 1:1 substitution ratio minimizes lattice distortion while maintaining computational tractability, consistent with previous DFT codoping studies [30]–[32]. Therefore, this configuration was adopted as a representative and physically meaningful model for evaluating the electronic and optical properties of N/Y-codoped TiO₂.

The first-principles calculations in this work were performed from DFT as implemented in the Cambridge Serial Total Energy Package (CASTEP) computer code. The structural, electronic, and optical properties of N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped anatase TiO₂ were calculated using the best exchange correlation functional, which is from the generalized gradient approximation in the Perdew-Burke-Ernzerhof for solids (GGA-PBEsol) [33]. The cut-off energy of 380 eV was required for N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped anatase TiO₂, while the *k*-point sampling of 4×4×2, 4×4×2, and 3×3×2 was set in Monkhorst-Pack grid for the N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂, respectively, which is within the convergence value as verified from the convergence test. Pseudo atomic calculation were 2*s*² 2*p*⁴ for O, 3*s*² 3*p*⁶ 3*d*² 4*s*² for Ti, 2*s*² 2*p*³ for N and 3*s*² 3*p*⁶ 3*d*³ 4*s*² 4*d*¹ for Y. The geometrical optimization was set at 5.0×10⁻⁶ eV/atom for total energy, 0.01 eV/Å for maximum force, 0.02 GPa for maximum stress, and 5.0×10⁻⁴ Å for maximum displacement.

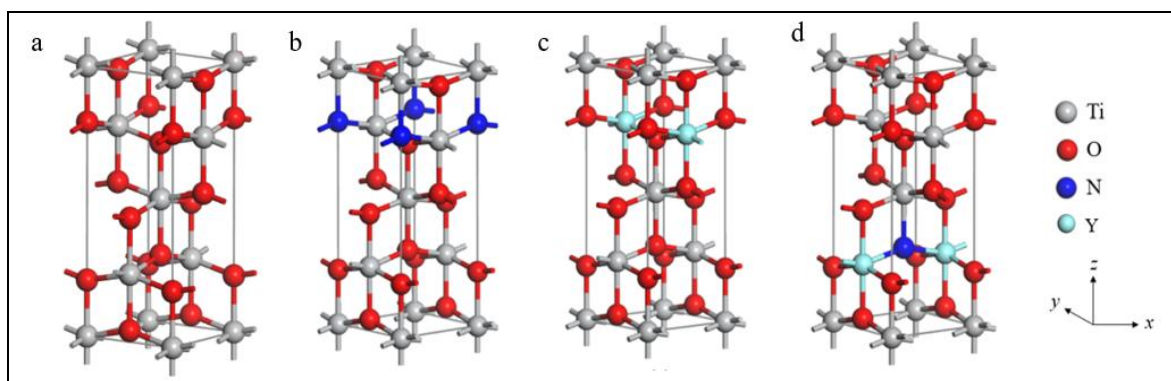


Figure 1. Crystal structures of (a) pure TiO₂; (b) N-doped TiO₂; (c) Y-doped TiO₂; (d) N/Y-codoped TiO₂

Results and Discussion

Structural parameters

The optimized lattice parameters of pure TiO₂, N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ are all listed in **Table 1**. The lattice constants of pure TiO₂ are consistent with those in previous experimental work [34]. This result implies that our calculating methods are acceptable and that the calculated outcomes are reliable.

The optimized cell characteristics for pure TiO₂ and TiO₂ that has been doped with N and Y are compared to examine the changes in geometrical structure brought on by the non-metal and metal dopants. It is found that the lattice parameters of N-doped TiO₂ were overestimated by 2.66% and 1.37% for *a* and *b* parameters, respectively. While lattice parameters of Y-doped TiO₂ were overestimated by 4.03% (*a* = *b* parameters) over pure TiO₂. For the *c* parameter, the underestimation of N-doped TiO₂ compared to pure TiO₂ is by 0.27% and for Y-doped TiO₂, the overestimation is by 12.01% compared to pure TiO₂. Compared with the lattice constant ratio (*c/a*), the underestimation and overestimation of N-doped TiO₂ and Y-doped TiO₂ are by 2.87% and 10.85% over pure TiO₂. For the volume, the N-doped TiO₂ and Y-doped TiO₂ expand their volume by 3.75% and 21.21%, respectively. For N/Y-codoped TiO₂, it is found that lattice parameters (*a*, *b*, and *c*), lattice constant ratio (*c/a*), and volume (*V*) of the codoping are larger than pure TiO₂. N/Y-codoped TiO₂ was overestimated by 2.55%, 3.36% and 15.35% for *a*, *b*, and *c*, respectively, over pure TiO₂. Compared with the lattice constant ratio (*c/a*) of pure TiO₂, the overestimation of N/Y-codoped TiO₂ is 12.48%. For the volume of N/Y-codoped TiO₂, their unit cell volume expands by 22.23%. The increased cell volume of TiO₂ suggests a larger surface area and higher dye-loading capability. This is because the atomic radius of Y (2.12 Å) is larger than Ti (1.76 Å). This optimized TiO₂ structure was used for further calculations to gain a deeper theoretical understanding of its properties.

The average bond length in pure TiO₂, N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ for Ti-O, O-O, N-Ti, V-O, and Y-O are listed in Table 2. In N-doped TiO₂, the Ti-N (1.967 Å) bond lengths are closely similar to Ti-O (1.959 Å) in pure TiO₂ [35]. However, in Y-doped TiO₂, the Y-O (2.511 Å) is longer than Ti-O (1.961 Å) by +0.10%, suggesting that the Y-O bond has weaker covalency than the Ti-O bond. Generally, O-O bond lengths are longer compared to Ti-O, N-Ti, and V-O. The Ti-O bond lengths slightly elongated after the N/Y-codoped TiO₂, as compared to the Ti-O bond lengths before codoping. However, the Y-O is longer than Ti-O and O-O bond lengths. The Y-O bond (2.662 Å) is longer than the Ti-O bond (1.959 Å) by approximately 0.51% in N/Y-codoped TiO₂. It should be noted that the longer bond length leads to a weaker bond strength. Generally, the increase of bond length between cation and anion may decrease the coupling constant and consequently decrease the band gap according to spectroscopic analysis of a set of tetrahedrally coordinated semiconductors [36].

Band structures

The incorporation of N and Y atoms in anatase TiO₂ not only influences the crystal structures but also affects its electronic properties. The band structure of pure TiO₂ [37] was calculated using the first-principles method and is shown in **Figure 2(a)**. The theoretical band gap of anatase TiO₂ calculated using the GGA-PBEsol functional is 2.149 eV, which is lower than the experimental value of approximately 3.2 eV, which is in agreement with previous DFT studies [38–40]. This underestimation is a well-known limitation of conventional GGA-based DFT methods, as they typically neglect the quasiparticle self-energy correction. Despite this, the relative differences between the doped and codoped systems remain valid for understanding the qualitative effects of N and Y incorporation on the electronic structure. The band structures of N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped anatase TiO₂ along the high symmetry directions of the Brillouin zone (BZ) at G(0,0,0)-F(0,0.5,0)-Q(0,0.5,0.5)-Z(0,0,0.5)-G(0,0,0) are shown in **Figure 2(b)-2(d)**.

Table 1. Lattice constants (a , b , c) and volumes (V) of pure TiO₂, N-doped TiO₂, Y-doped TiO₂ and N/Y-codoped TiO₂ by geometry optimization

	Present works				Experimental [34]
	Pure TiO ₂	N-doped TiO ₂	Y-doped TiO ₂	N/Y-codoped TiO ₂	Pure TiO ₂
a (Å)	3.720	3.819 (+2.66%)	3.870 (+4.03%)	3.815 (+2.55%)	3.787
b (Å)	3.720	3.771 (+1.37%)	3.870 (+4.03%)	3.845 (+3.36%)	3.787
c (Å)	9.600	9.574 (-0.27%)	10.753 (+12.01%)	11.074 (+15.35%)	9.514
c/a (Å)	2.581	2.507 (-2.87%)	2.779 (+10.85)	2.903 (+12.48%)	2.512
V (Å ³)	132.88	137.86 (+3.75%)	161.07 (+21.21%)	162.42 (+22.23%)	136.44

(+) overestimate; (-) underestimate compared to pure TiO₂

Table 2. Average bond length of doped TiO₂ after geometry optimization

	Pure TiO ₂	N-doped TiO ₂	Y-doped TiO ₂	N/Y-codoped TiO ₂
Ti-O (Å)	1.959	1.970 (+0.56%)	1.961 (+0.10%)	1.969 (+0.51%)
O-O (Å)	2.626	2.637 (+0.42%)	2.775 (+5.67%)	2.769 (+5.45%)
N-Ti (Å)	–	1.967	–	1.926
Y-O (Å)	–	–	2.511	2.662

(+) overestimate compared to pure TiO₂

When the non-metal and metal elements are doped in TiO₂, the energy gaps are reduced. For pure anatase TiO₂, the energy band gap is 2.149 eV, while the N-doped TiO₂ and Y-doped anatase TiO₂ show a reduced band gap to 1.534 eV and 2.142 eV, respectively. The N doping causes a slight decrease in the band gap to 1.534 eV in connection with the formation of localized N $2p$ states just above the valence band maximum (VBM) of TiO₂, which can be seen in the density of states (DOS). A decrease in conduction band minimum (CBM) and VBM causes the band gap to widen due to intense nitrogen doping is similar to that described by Liu and co-workers [41]. For Y-doped TiO₂, the CBM moves slightly downward and reduces the band gap by 0.33% compared to pure TiO₂. In addition, the Fermi level moves into the VB and shows the property of a p-type semiconductor. The reduction of energy gaps from the TiO₂ codoping is due to the lowering of the CB. The N/Y-codoped anatase TiO₂ shows a reduced band gap to 1.880 eV. The slight decrease in bandgap values for the rare earth doped could be attributed to size quantization or charge carrier spatial confinement. This is due to the fact that the electron-hole pairs are now considerably closer together, and the Coulombic interaction between them can no longer be ignored, resulting in a larger overall kinetic energy [42].

Density of states

The partial density of states (PDOS) and total density of states (DOS) of N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ are given in **Figure 3** to explore the electronic structures. It is well known that the conduction band of TiO₂ is composed of Ti $3d$ orbital, and the valence band is composed of O $2p$ orbital and in agreement with previous work by other researchers [43]. For the N-doped TiO₂, the orbital hybridization of Ti $3d$, O $2p$, and N $2p$ forms the valence band of N-doped TiO₂. After N doping, it is very clear that the empty band appeared in the middle of the forbidden band of TiO₂. The N $2p$ orbital lies near the CBM of TiO₂, which decreases the band gap of TiO₂ and improves the photocatalytic activity.

From the DOS of Y-doped TiO₂, it can be seen that d states mostly contribute to the conduction zone, while p states contribute mainly to the valence zone. The obtained results are also in agreement with what has been established for Y atoms, and that is the fact that their presence in the structure narrows the band gap, which is of particular importance for photocatalysis [44, 45]. In the case of the N/Y-codoped TiO₂, the Y $4d$ state dominates the conduction band, which is located close to the bottom of the conduction band, while the N $2p$ state dominates the valence band, which is located close to the top of the valence band. From the DOS of N/Y-codoped TiO₂, it can be concluded that the absorption edge transition occurs between N $2p$ and Y $4d$ states.

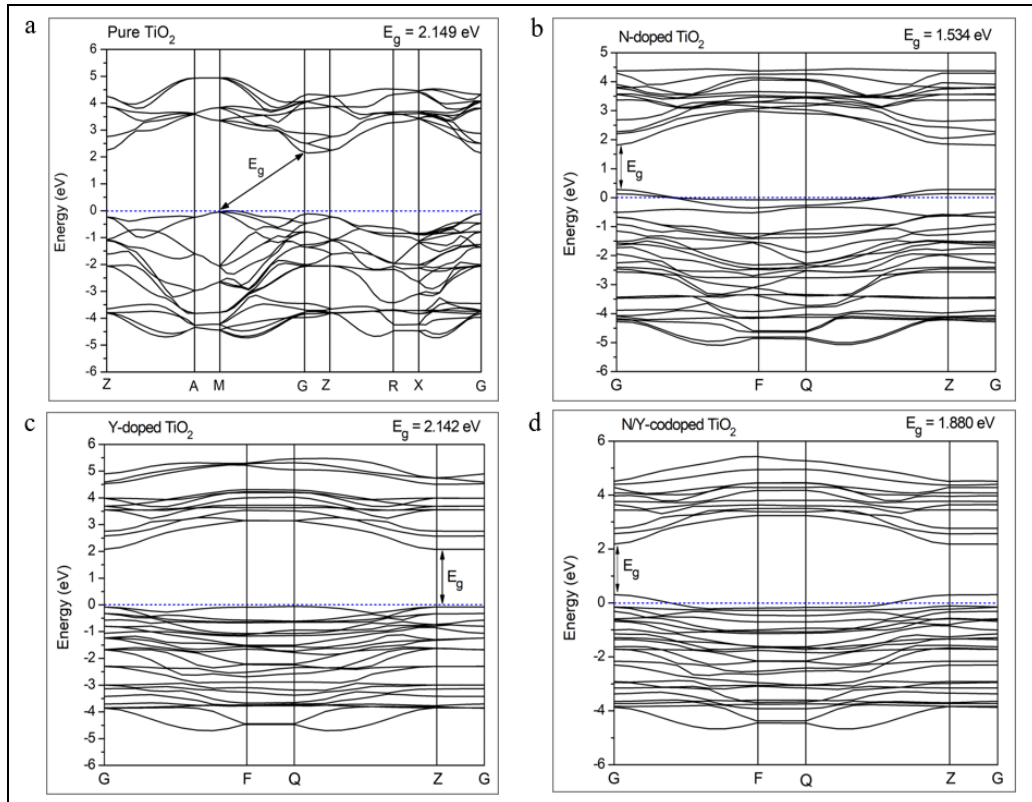


Figure 2. The band structures of (a) pure TiO_2 ; (b) N-doped TiO_2 ; (c) Y-doped TiO_2 ; (d) N/Y-codoped TiO_2 . The Fermi energy is set to zero as indicated by the blue dashed line. The calculated values are smaller than experimental results due to the well-known underestimation of band gaps by the GGA-PBESol functional

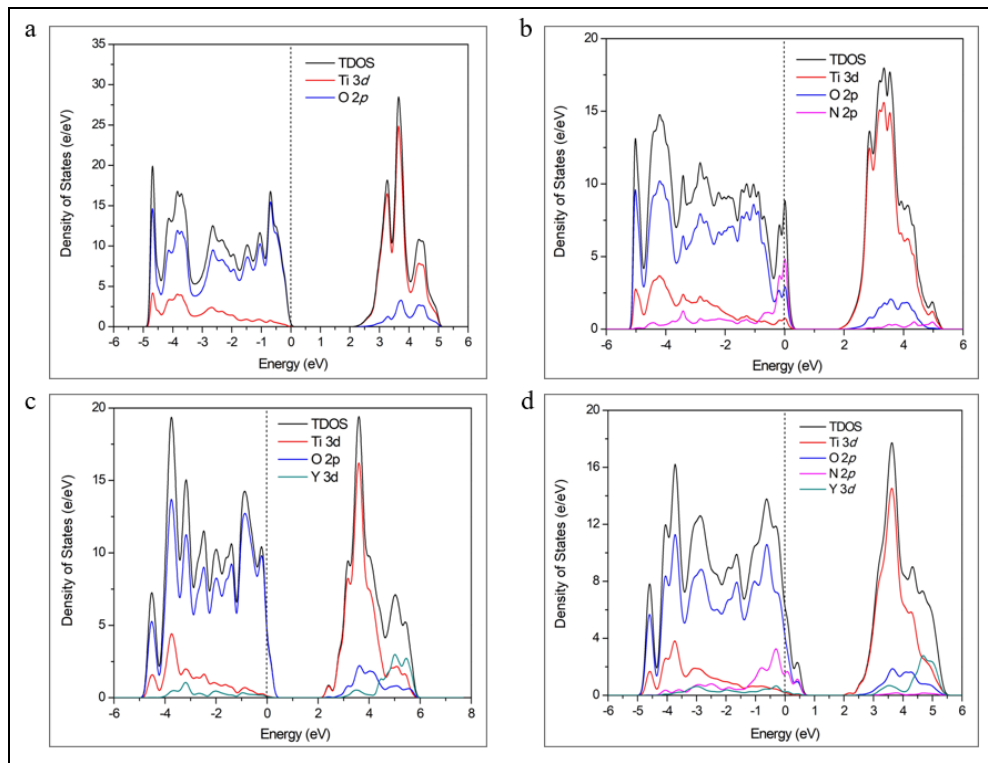


Figure 3. The TDOS and corresponding PDOS of (a) pure TiO_2 , (b) N-doped TiO_2 , (c) Y-doped TiO_2 , and (d) N/Y-codoped TiO_2 . The Fermi level is indicated by the black dashed line

Absorption spectra

The absorption coefficients of pure TiO₂, N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ against photon energy and wavelength are displayed in **Figure 4**. The pure TiO₂ exhibits an absorption edge at approximately 2.149 eV (~577 nm), consistent with its ability to primarily absorb UV light. After doping, a clear red-shift in the absorption edge is observed. The absorption onsets shift to ~1.534 eV (~809 nm) for N-doped TiO₂, ~2.142 eV (~579 nm) for Y-doped TiO₂, and ~1.880 eV (~660 nm) for N/Y-codoped TiO₂. This behavior corresponds directly to the calculated band gap values and confirms that doping introduces impurity states that narrow the band gap. Among all systems, the N/Y-codoped TiO₂ exhibits the greatest visible-light absorption enhancement, indicating a synergistic effect where N contributes to valence band elevation while Y influences conduction band states, thereby facilitating improved charge carrier excitation under visible illumination. In conclusion, the electrons in the valence band can be excited to the conduction band by absorption of the visible light, and non-metal and metal doping are ideal to effectively utilize solar energy in the visible light region.

For N-doped TiO₂, the photo-excited carrier will be trapped by the N 2*p* state and then transferred into the conduction band. As shown in **Figure 4 (b)**, the Y-doped TiO₂ system does not show good absorption activity in the visible light region. For the N/Y-codoped TiO₂ system, the transition energy becomes smaller due to the split of the impure state, rendering a more obvious red-shift. Thus, the photocatalytic efficiency is enhanced by suppressing the recombination of electron-hole pairs. Compared with the N-doped TiO₂ and Y-doped TiO₂, the N/Y-codoped TiO₂ is expected to be a more active photocatalyst.

Conclusion

In this study, the structural, electronic, and optical properties of N-doped, Y-doped, and N/Y-codoped anatase TiO₂ were investigated using first-principles calculations within the GGA-PBEsol framework. The results indicate that N doping mainly contributes to the upward shift of the valence band through the introduction of N 2*p* states, while Y doping modifies the conduction band through Y 4*d* states. When codoped, these effects act synergistically, resulting in a more pronounced reduction in the band gap (1.880 eV) compared to single-doped systems and promoting an enhanced red-shift in the optical absorption edge toward the visible-light region. This demonstrates that N/Y codoping is a promising strategy for improving visible-light-driven photocatalytic activity of TiO₂.

However, the present work is limited to bulk electronic structure calculations and does not explicitly consider surface states, charge recombination dynamics, or adsorption/reactive intermediates, which are critical to actual photocatalytic performance. Therefore, future studies should include experimental synthesis and photocatalytic testing, as well as surface reactivity and carrier transport analyses, to verify and further optimize the functional role of N/Y codoping. These additional investigations may help establish a stronger correlation between theoretical predictions and real-world photocatalytic applications.

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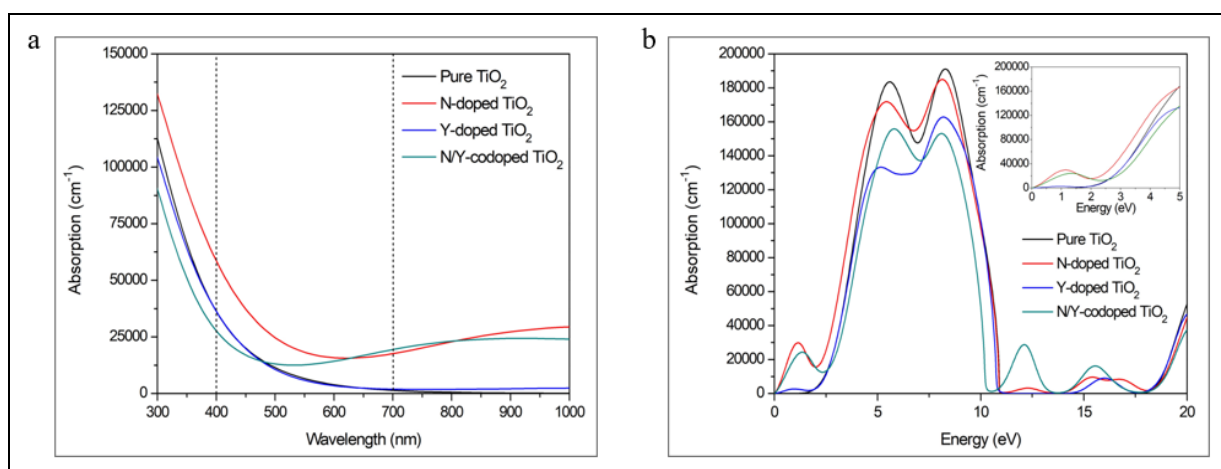


Figure 4. The optical absorption spectra of pure TiO₂, N-doped TiO₂, Y-doped TiO₂, and N/Y-codoped TiO₂ against (a) wavelength (nm) and (b) energy (eV)

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